Correlation functions

radial distribution function (RDF, also pair correlation/distribution function), g(r) = probability of finding a particle at distance r (from another particle), normalized to ideal gas

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- structure factor (diffraction \rightarrow Fourier transform of g(r))
- angular correlation function good for small nonspherical molecules
- time autocorrelation functions

Order parameters

- Ordering in the *z*-direction: $\langle \sum_i \cos \theta_i \rangle$: \uparrow vs. \downarrow
- For symmetric molecules: $\left\langle \sum_{i} \left(\frac{3}{2} \cos^2 \theta_i \frac{1}{2} \right) \right\rangle$: \ddagger vs. \leftrightarrow
- Nematic liquid crystal the "director" is not known:

$$Q_{\alpha\beta} = \left\langle \sum_{i} \left(\frac{3}{2} \hat{n}_{\alpha,i} \hat{n}_{\beta,i} - \frac{1}{2} \delta_{\alpha\beta} \right) \right\rangle, \text{ eigenvector of the max. eigenvalue} = \text{director}$$

Formulas detecting local order (e.g., tetrahedral around water molecules as Steinhardt's parameters based on spherical harmonics), onset of crystallization can be detected ...

Structure of fluids – correlation functions





randomly distributed molecules (ideal gas)

liquid

g(r) = pair correlation function = radial distribution function = probability density of finding a particle r apart from another particle, normalized so that for randomly distributed particles (ideal gas) it is 1

Structure of fluids – correlation functions



How to obtain structure – experiment



Diffraction (neutrons, X-ray, electrons) \Rightarrow "structure factor"

inverse Fourier transform \Rightarrow RDF

How to obtain structure



Argon, hard spheres, water





- The structure of simple fluid (argon, HS) is organized by shells.
- The structure of water is determined by the tetrahedral geometry of hydrogen bonds.
- After several molecular diameters, the correlations decay to zero.

Running coordination number

Also "cumulative radial distribution function"

 $\rho = N/V =$ number density

$$N(r) = 4\pi\rho \int_0^r g(r')r'^2 dr'$$

For r_{min} = first minimum on the RDF curve, $N(r_{min})$ = "coordination number" = averaged number of molecules in the first shell



RDF from simulations – simple

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Histogram of the number of particle pairs, \mathcal{N}_i , so that

 $r \in [r_i - \Delta r/2, r_i + \Delta r/2)$ alternatively: $\mathcal{I}_i = [r_i, r_i + \Delta r)$

The volume of the shell

$$\Delta V_i = \frac{4\pi}{3} \left[\left(r_i + \frac{\Delta r}{2} \right)^3 - \left(r_i - \frac{\Delta r}{2} \right)^3 \right]$$



Mean number of molecules around a selected particle in case of uniformly distributed molecules (ideal gas, $\rho = N/V$):

 $\rho \Delta V_i$

Sum over all particles (1/2 to count each pair just once):

$$\mathcal{N}_i^{\text{id.gas}} = \frac{N}{2} \rho \Delta V_i$$

Radial distribution function:

$$g(r_i) = \frac{\langle \mathcal{N}_i \rangle}{\mathcal{N}_i^{\text{id.gas}}} = \frac{2 \langle \mathcal{N}_i \rangle}{N \rho \Delta V_i}$$

Another example – brine

RDFs of supersaturated (10.5 mol/kg) solution of NaCl:



Another example – brine

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CNs of supersaturated (10.5 mol/kg) solution of NaCl:



- http://old.vscht.cz/fch/software/simolant or Google simolant
- Download simolant-win32.zip
- Create a folder and unpack SIMOLANT there.
 Do not run directly from simolant-win32.zip!
- Run simolant.exe
- **Hint:** The calculated data are exported to file simolant.txt with a decimal point. If you like decimal comma (useful with Czech localization), click **I** comma in panel "Expert".
- **Hint:** If you restart SIMOLANT, the old simolant.txt is renamed to simolant.bak. The export name simolant can be changed by Menu: File \rightarrow Protocol name..

Simolant: observe RDF by yourself

- Menu: <u>Prepare system</u> \rightarrow <u>Periodic liquid</u>
- For smoother results, move slider "measurement block" to maximum
- Increase/decrease temperature a bit (slider "T"). Try a different pressure (slider "P"). NB: the RDF graph is scaled from 0 to maximum!

Set the temperature to maximum (T = 5): only a small first peak should remain.

- Decrease temperature below 0.4 to crystallize.
- To obtain a nearly perfect crystal without defects, select Menu: Prepare system \rightarrow Periodic crystal. Again, you may try to heat/cool. You should get: $r_1 \quad r_2r_3$





• Apparently $r_3/r_1 = 2$ because r_3 is the second neighbor •••

Similarly,
$$r_2/r_1 = \sqrt{3}$$



Radial distribution function in the *NVT* **ensemble**

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3D (e.g., in periodic b.c.), good for a homogeneous and isotropic liquid:

$$g(r) \equiv g(r_{12}) = \frac{N(N-1)}{\rho^2 Q_{\text{NVT}}} \int \dots \int \exp[-\beta U(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N)] \, \mathrm{d}\vec{r}_3 \dots \, \mathrm{d}\vec{r}_N$$

Equivalently

$$g(r) = \left(1 - \frac{1}{N}\right) V \left< \delta(\vec{r}_{12} - \vec{r}) \right>$$

For a mixture:

$$g_{ij}(r) = V \left< \delta(\vec{r}_{12} - \vec{r}) \right>$$

Normalization (fluid):

$$\lim_{N\to\infty,r\to\infty}g(r)=1$$

NB: ideal gas at finite N: g(r) = 1 - 1/N (e.g., in periodic b.c.)

Number of particles around one chosen particle (in NVT):

$$\int_V \rho g(\vec{r}) \mathrm{d}\vec{r} = N-1$$

 $g(r_{12}) = \frac{\rho_2(\vec{r}_1, \vec{r}_2)}{\rho_1(\vec{r}_1)\rho_2(\vec{r}_2)},$ where ρ_n is the *n*-particle distribution function, $\rho_1 =$

 ρ in a homogeneous fluid.

Or (mixtures, in a region):

Calculation of RDF in simulations – exact

Histogram of the count of pairs of particles, \mathcal{N}_i , so that $r \in \mathcal{I}_i$

$$\begin{split} \mathcal{I}_{i} &= [r_{i} - \Delta r/2, r_{i} + \Delta r/2), \text{ optionally } \mathcal{I}_{i} = [r_{i}, r_{i} + \Delta r) \\ r_{i} &= i\Delta r, i = 1, \dots, i_{\max} \\ \mathcal{N}_{i} \rangle &= \frac{1}{Q_{\mathsf{NVT}}} \sum_{j < k} \int_{r_{jk} \in \mathcal{I}_{i}} \exp[-\beta U(\vec{r}^{N})] d\vec{r}^{N} \\ &= \frac{1}{Q_{\mathsf{NVT}}} \binom{N}{2} V \int_{r_{12} \in \mathcal{I}_{i}} \left\{ \int \exp[-\beta U(\vec{r}^{N})] d\vec{r}_{3} \dots d\vec{r}_{N} \right\} d\vec{r}_{12} \\ &= \frac{N}{2} \rho \int_{\mathcal{I}_{i}} g(r) d\vec{r} \\ &\approx \frac{N^{2}}{2V} g(r_{i}) \Delta V_{i} \end{split}$$

The formula again:

$$g(r_i) = \frac{2\langle \mathcal{N}_i \rangle}{N \rho \Delta V_i}$$

Correlation function and thermodynamics



For simple fluid (point particles interacting by a pair spherically symmetric potential): **Residual internal energy:**

$$\langle U \rangle = \frac{1}{Q_{\text{NVT}}} \int \sum_{i < j} u_{ij}(r_{ij}) e^{-\beta U} d\vec{r}_1 \dots d\vec{r}_N$$

$$\langle U \rangle = \binom{N}{2} \frac{V}{Q_{\text{NVT}}} \int e^{-\beta U} 4\pi r_{12}^2 dr_{12} u(r_{12}) d\vec{r}_3 \dots d\vec{r}_N$$

$$\langle U \rangle = \frac{N}{2} \rho \int u(r) g(r) d\vec{r} = 2N\pi\rho \int u(r) g(r) r^2 dr$$

Pressure:

$$\frac{\beta P}{\rho} = 1 - \frac{2\pi}{3}\beta\rho \int g(r)u'(r)r^3 dr$$





- Distance of atom \vec{r}_j from plane $\vec{v} \cdot \vec{r} = 0$ is $\vec{v} \cdot \vec{r}_j$ (for $|\vec{v}| = 1$).
- Distance from wave source (wave plane of eq. $\vec{v} \cdot \vec{r} = \text{const}$) via atom r_j to the detector ($\vec{v'} \cdot \vec{r} = \text{const}$) is $\vec{v} \cdot \vec{r}_j \vec{v'} \cdot \vec{r}_j + \text{const}$ Definition $k = 2\pi/\lambda$ is common in physics
- Wave vector $\vec{k}_{inc} = k_{inc} \vec{\nu}$, $k_{inc} = 2\pi/\lambda$
- Formally the incident wave is (but phase)

Definition $k = 2\pi/\lambda$ is common in physics wheras $k = 1/\lambda$ in crystallography; then, factor 2π is in the exponential: exp[$2\pi i \vec{k} \cdot \vec{r}$].

$$\exp[i(k_{\text{inc}}\vec{\nu} - k_{\text{inc}}\vec{\nu}') \cdot \vec{r}_j] = \exp[i\vec{k} \cdot \vec{r}_j]$$

where

$$\vec{k} = k_{\text{inc}}\vec{\nu} - k_{\text{inc}}\vec{\nu}', \qquad k = |\vec{k}| = 2\sin(\alpha/2)k_{\text{inc}} \approx \alpha k_{\text{inc}}$$

Structure factor

Waves scattered by all atoms interfere:

$$K = \sum_{j=1}^{N} \exp[i(k_{\text{inc}}\vec{\nu} - k_{\text{inc}}\vec{\nu}') \cdot \vec{r}_j] = \sum_{j=1}^{N} \exp[i\vec{k} \cdot \vec{r}_j], \quad \text{signal} \propto \langle |K|^2 \rangle$$

Definition of the structure factor (pure compound):

$$S(\vec{k}) = \left\langle \frac{1}{N} \left| \sum_{j=1}^{N} \exp(-i\vec{k} \cdot \vec{r}_j) \right|^2 \right\rangle$$

It is a function of vector \vec{k} , we may imagine it in the periodic b.c. in a cube of edge L:

$$\vec{k} = \frac{2\pi\vec{n}}{L}, \ \vec{n} \in \mathbb{Z}^3$$

Similarly, a 3D RDF is (not exactly isotropic even for a liquid)

$$g(\vec{r}_{12}) = \frac{N(N-1)}{\rho^2 Q_{\text{NVT}}} \int_{L^3} \dots \int_{L^3} \exp[-\beta U(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N)] \, \mathrm{d}\vec{r}_3 \dots \, \mathrm{d}\vec{r}_N$$

Structure factor and RDF

In a periodic box
$$V = L^3$$
; $\vec{k} = 2\pi\vec{n}/L$:

$$S(\vec{k}) = \left\langle \frac{1}{N} \left| \sum_{j=1}^{N} \exp(-i\vec{k} \cdot \vec{r}_j) \right|^2 \right\rangle$$

$$= \left\langle \frac{1}{N} \sum_{j=1}^{N} \sum_{l=1}^{N} \exp(-i\vec{k} \cdot \vec{r}_l) \exp(i\vec{k} \cdot \vec{r}_l) \right\rangle = \left\langle \frac{1}{N} \sum_{j=1}^{N} \sum_{l=1}^{N} \exp(i\vec{k} \cdot \vec{r}_{jl}) \right\rangle$$

$$= 1 + \frac{1}{N} \left\langle \sum_{j

$$= 1 + \rho \int_{I^3} g(\vec{r}) \cos(\vec{k} \cdot \vec{r}) d\vec{r}$$$$

The integral diverges as $V \to \infty \Rightarrow$ we subtract $0 = \int_V \cos(\vec{k} \cdot \vec{r}) d\vec{r}$, $\vec{k} = 2\pi \vec{n}/L$ (other method: convergence factor $\exp(-\alpha k^2)$, $\lim_{\alpha \to 0}$)

$$S(\vec{k}) = 1 + \rho \int_{V} [g(\vec{r}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{r} \stackrel{V \to \infty}{=} 1 + \rho \int [g(\vec{r}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{r}$$

Structure factor and RDF



Since $g(\vec{r}) = g(-\vec{r})$, we have

$$S(\vec{k}) - 1 = \rho \int [g(\vec{r}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{r} = \rho \int [g(\vec{r}) - 1] e^{-i\vec{k} \cdot \vec{r}} d\vec{r}$$

and the 3D inverse Fourier transform is

$$g(\vec{k}) - 1 = \frac{1}{8\pi^3 \rho} \int [S(\vec{k}) - 1] e^{i\vec{k}\cdot\vec{r}} d\vec{k}$$

Once again using $S(\vec{k}) = S(-\vec{k})$:

$$g(\vec{k}) - 1 = \frac{1}{8\pi^3 \rho} \int [S(\vec{k}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{k}$$

The $k \rightarrow 0$ limit:

NVT: **S(0)** = 0

$$\mu VT$$
: **S(0)** = 1 + $\rho \int [g(\vec{r}) - 1] d\vec{r} = k_{\rm B} T \left(\frac{\partial \rho}{\partial \rho}\right)_T$ (compressibility equation)

Isotropic structure factor

Let's choose $\vec{k} = (0, 0, k)$ and spherical coordinates (ϕ, θ, r) , then $\vec{k} \cdot \vec{r} = kr \cos \theta$ and (g and S are odd functions)

$$S(k) - 1 = \rho \int [g(\vec{r}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{r} \quad (\text{subst. } z = \cos\theta)$$
$$= \rho \int_0^\infty r^2 dr \int_0^{2\pi} d\phi \int_{-1}^1 [g(r) - 1] \cos(krz) dz$$
$$= \frac{4\pi\rho}{k} \int_0^\infty r[g(r) - 1] \sin(kr) dr$$

In the same way we can evaluate the inverse Fourier transform

$$g(\vec{k}) - 1 = \frac{1}{8\pi^{3}\rho} \int [S(\vec{k}) - 1] \cos(\vec{k} \cdot \vec{r}) d\vec{k}$$
$$= \frac{1}{2\pi^{2}r\rho} \int_{0}^{\infty} k[S(k) - 1] \sin(kr) dk$$

S(k) from simulations

 $S(\vec{k})$ in simulations can be calculated directly by the same code as the k-space part of the Ewald summation. If needed, we may sphericalize it:

$$S(k) = \sum_{k=|\vec{k}|} S(\vec{k}) / \sum_{k=|\vec{k}|} 1$$

S(k) can be obtained by the inverse Fourier transform of g(r). The RDF g(r) must have a long range and/or it has to be filled (by 1) for long distances.

If one wishes, 3D $S(\vec{k})$ can be obtained by the inverse Fourier transform from RDF $g(\vec{r})$ and vice versa

- Simulation-based *S*(*k*) for a mixture (where individual site-site RDFs need not be available) allows for a direct comparison with the experiment
- \bigcirc S(k) and S(\vec{k}) may detect unwanted crystallization of a supercooled system

Structure factor



Structure factor for multiatomic systems

 $+\frac{23/29}{s08/4}$

Structure factor for a "mixture" of atoms (b_i = coherent scattering length)

$$S(k) = 1 + N \frac{\langle |Q(\vec{k})|^2 \rangle - \sum_j b_j^2}{\left(\sum_j b_j\right)^2}$$

$$Q(\vec{k}) = \sum_{j} b_{j} \exp[-2\pi i \vec{k} \cdot \vec{r}_{j}/L]$$

$$S = \sum_{I} \sum_{J} w_{IJ} S_{IJ}, \quad w_{IJ} = \frac{N_I b_I N_J b_J}{(\sum_{I} N_I b_I)^2}$$

 N_I = number of atoms of type I ($\sum_I N_I b_I = \sum_j b_j$)

Reverse Monte Carlo



Input: experimental RDF

Output: *u*(*r*) so that the RDF is best reproduced

Not unique – other conditions on u(r) needed

Exercise: structure factor of NaCl crystal/melt



Scattering lenghts [fm]: Na: 3.63, Cl: 9.577 (³⁵Cl: 11.65, ³⁷Cl: 3.08)

Tasks:

- Simulate a crystal at 1000 K
- Switch to NVT and generate a trajectory (about 100 configurations)
- Calculate the SF directly
- Compare with inverted RDF
- Melt and repeat for the liquid
- Try different isotopes

You need: nacl.ble nacl.def nacl.cfg na.mol na.gol cl.mol cl.gol

```
Make your folder, and install (to your folder) by:
guest@403-a324-01:~/VY$ mkdir YOU
guest@403-a324-01:~/VY$ cd YOU
guest@403-a324-01:~/VY$ unzip /home/jiri/nacl.zip
```

```
Inspect carefully nacl.def!
```

Exercise: structure factor of the NaCl crystal

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A crystal is prepared in nacl.def Start + equilibration

start simulation; all variables are default: guest@403-a324-01:~/VY\$ cook nacl -s init="start" start from prepared configuration ; end of data set + start of simulation

after the run has finished: key="cp" show the convergence profile key="show" show the trajectory ; another 100 steps if not equilibrated

final run:

tau.P=0 use NVT now

init="start" use the configuration, measurements + trajectory restarted

- ; end of data set + start of simulation
- key="rdf" calculate + show RDF
- key="cn" show the coordination numbers
- key="quit" or ctrl-d = control returns to the command prompt

Exercise: structure factor of the crystal directly

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The structure factor will be calculated from the trajectory, nacl.plb, recorded in previously.

- b data: optionally change b_i at the end of table Lennard-Jones in nacl.ble
- How long is the trajectory? guest@403-a324-01:~/VY\$ plbinfo nacl.plb
- Calculate the structure factor (-m1 = reads nacl.plb) guest@403-a324-01:~/VY\$ cook nacl -m1 -s reread.from=1 reread.to=101 read nacl.plb from .. to tau.P=0 NVT (do not change the box) el.kappa=15/2/pi ; max.k-vector
- The result is in nacl.sfr. To show: guest@403-a324-01:~/VY\$ plot nacl.sfr

Exercise: structure factor of the NaCl crystal



Now the same by the Fourier transform of the RDF

Calculate the RDF first:

```
guest@403-a324-01:~/VY$ rdfg nacl -g -p
```

Fourier transform for a mixture with given b_i : guest@403-a324-01:~/VY\$

cat nacl.*.g | sfourier 100 30 216 NA:3.63 CL:9.577 >| nacl.sfg

```
Comparison:
guest@403-a324-01:~/VY$ plot nacl.sfr nacl.sfg
nacl.sfr = direct result
nacl.sfg = RDF-inverted result
```

Exercise: structure factor of the NaCl melt

We will melt the sample, equilibrate, and simulate the melt.

```
Melting:
guest@403-a324-01:~/VY$ cook nacl -s
T=5000 a lot at start, to melt the crystal
tau.P=0 but NVT, so that it does not evaporate
no=50 ; start 50 steps
```

Cooling the melt

T=1400 ; temperature of the melt and cooling in NVT

Equilibrating of the melt sample at given temperature, pressure bulkmodulus=2e9 approx. bulk modulus of the liquid for NPT [in Pa] tau.P=1 time constant of the barostat no=100; running at selected temperature, pressure in NPT key="cp" to have a look whether the system has equilibrated ; repeat until equilibrated

Final run tau.P=0 NVT again to calculate RDF + write a trajectory init="start";

To be followed by the same algorithm as in the case of crystal...