Vectors

Kindergarten: vector = $(v_1, ..., v_n)$, $v_i \in \mathbb{R}$

Quantum kindergarten: $v_i \in \mathbb{C}$

Mathematics: **vector space** (linear space) is defined by the axioms:

For vectors u, v, w and numbers $a, b \in \mathbb{R}$ or \mathbb{C} (a field in general):

$$u+(v+w) = (u+v)+w$$

$$u+v = v+u$$

$$\exists \text{ null vector } 0:v+0 = v$$

$$\exists \text{ opposite vector } -v:v+(-v) = 0$$

$$a(bv) = (ab)v$$

$$1v = v$$

$$a(u+v) = au+av$$

$$(a+b)v = av+bv$$

Notation: v, \mathbf{v} , \vec{v} (real in 2D, 3D), v, $|v\rangle$ ("ket"), v_i (?)

Orthogonal and orthonormal bases

Orthogonal basis = all vectors are perpendicular.

Orthonormal basis = also normalized.

$$b^{(i)} \cdot b^{(j)} = \delta_{ii}$$

Components of ν in an orthonormal basis:

$$v_i = v \cdot b^{(i)} \ \Rightarrow \ v = \sum v_i b^{(i)} = (v_1, \dots, v_n)_b$$

Scalar product:

$$u\cdot v=\sum u_iv_i$$

Scalar product in $\mathbb C$ in physics

$$\langle u|\nu\rangle = \sum u_i^* v_i$$

Linear dependence

A set of nonzero vectors $v^{(i)}$, i = 1..m, is **linearly dependent** if there is a null linear combination with at least one of a_i nonzero:

$$\sum a_i v^{(i)} = 0$$

A linearly independent set of vectors such that any vector (of given space) can be expressed as its linear combination is called a basis

$$v = \sum_{i} v_i b^{(i)}$$

 $(1, 2, 3, 4)^{T}, (1, -2, 3, -4)^{T}, (1, 0, 1, 0)^{T}$

Example. Are the following vectors in \mathbb{R}^4 linearly dependent?

Example. Are the following vectors in \mathbb{C}^2 linearly dependent?

$$\begin{pmatrix} i \\ 1 \end{pmatrix}, \begin{pmatrix} 1+i \\ 1-i \end{pmatrix}$$

Example. Consider a linear space of functions of $x \in [0, 2\pi]$ with basis $\{1, \cos(x), \cos(2x), \cos(x), \cos(2x), \cos(x), \cos(2x), \cos(x), \cos(x),$ cos(3x),... Can function $cos^2(x)$ be expressed in this basis? yes: {1/2, 0, 1/2, 0, 0, 0, 0, ...}

Gram-Schmidt orthogonalization

A general basis $b^{(i)}$ can be orthogonalized by the Gram-Schmidt algorithm:

$$b^{(1)} \ := \ b^{(1)}/|b^{(1)}|$$

$$b^{(2)} \ := \ b^{(2)} - \langle b^{(1)} | b^{(2)} \rangle b^{(1)}, \ b^{(2)} := b^{(2)} / |b^{(2)}|$$

$$b^{(3)} \ := \ b^{(3)} - \langle b^{(1)} | b^{(3)} \rangle b^{(1)} - \langle b^{(2)} | b^{(3)} \rangle b^{(2)}, \ b^{(3)} := b^{(3)} / |b^{(3)}|$$

":=" means "assign to" as in computer code.

Bases used in a Hilbert space are usually orthogonal or orthonormal

Example. Find all orthonormal bases
$$\{b^{(1)},b^{(2)}\}$$
 in \mathbb{C}^2 for $b^{(1)}=(1,i)/\sqrt{2}$ $(b_1^{(1)}=1,b_2^{(1)}=i)$

$$\left\langle \left(\begin{array}{c} 1\\i \end{array}\right) \middle| \left(\begin{array}{c} x\\y \end{array}\right) \right\rangle = x - iy \stackrel{!}{=} 0 \implies x = iy \implies b^{(2)} = \frac{c}{\sqrt{2}} \middle| \left(\begin{array}{c} i\\1 \end{array}\right) \right\rangle, |c| = 1$$

Remember complex conjugate in the dot product, $i^* = -i$

more examples: see mmpc1.mw

Another notation:

Scalar (inner, dot) product

We need a richer structure!

Kindergarten: scalar product $\vec{u} \cdot \vec{v} = \sum u_i v_i$

Mathematics: (u, v) is a number (real, complex) obeying axioms:

$$(u, v) = (v, u)^* (* = complex conjugate)$$

$$(au, v) = a^*(u, v) \text{ (in physics)} \Rightarrow (u, av) = a(u, v)$$

$$= a(u, v) \text{ (in mathematics)} \Rightarrow (u, av) = a^*(u, v)$$

$$(u+v, w) = (u, w) + (v, w)$$

$$(u, u) \geq 0$$

$$(u, u) = 0 \Rightarrow u = 0 \text{ (null vector)}$$

Notation: $u^{\mathsf{T}} \mathsf{v}, \, u^{\mathsf{T}} \cdot \mathsf{v}, \, u^{\dagger} \mathsf{v}, \overset{\mathbf{b}}{} (u, v), \, \langle u, v \rangle, \, \vec{u} \cdot \vec{v}, \, \underline{u} \cdot v, \, \langle u | v \rangle \text{ (bra-ket)}, \, u_i v^i \text{ (covector-vector)}$

· (real spaces) or | (complex spaces) = sum over a pair of indices

Definition: If (u, v) = 0, vectors u, v are **perpendicular** $(u, u)^{1/2} = |u| = ||u|| = \mathsf{norm}^{\mathsf{c}}$

 $^{\mathbf{b}}$ symbol † = transpose + complex conjugate = adjoint = Hermitean (Hermitian) conjugate

csimilar space with a norm only (and complete) = Banach space; under some conditions $(u, v) = (|u+v|^2 - |u-v|^2)/4$

Linear form f (linear operator) assigns a number $f(v) \in \mathbb{R}$ (or \mathbb{C}) to a vector

Axioms: for linear forms f, g, number a, and a vector v:

$$(f+g)(v) = f(v)+g(v)$$

$$f(av) = af(v)$$

For finite *n* one can write (In infinite-dimension spaces there may be continuity problems):

$$f(u) = \sum_{i=1}^{n} f_i u_i$$

Otherwise in Hilbert spaces linear form ≈ scalar product:

$$f(v) = \sum f_i v_i = (f^*, v)$$

Linear form in Euclidean spaces (in some context) = covector, dual vector, covariant vector ("normal" vector = contravariant vector)

 \bigcirc covector = row vector (transposed) f^{T} , inverse transformation if a basis changes

Scalar product then is: $f(u)=f^{\mathsf{T}}\cdot u=f^{\mathsf{T}}u=f^{i}u_{i}$ (Einstein summation convention). In complex Hilbert spaces $^{\mathsf{T}}\!\!\to\! ^{\dagger}$

(Cauchy-)Schwarz inequality

Dot-product in \mathbb{R}^n : $\vec{x} \cdot \vec{y} = |\vec{x}||\vec{y}|\cos\theta \equiv xy\cos\theta \le xy$. For nonzero a, b (zero cases are trivial):d

$$b_{\perp} = b - \frac{\langle a|b \rangle}{a^2} a \Rightarrow \langle a|b_{\perp} \rangle = \langle a|b - \frac{\langle a|b \rangle}{a^2} a \rangle = \langle a|b \rangle - \frac{\langle a|b \rangle}{a^2} \langle a|a \rangle = 0$$

 $b = \frac{\langle a|b\rangle}{a^2}a + b_{\perp}$

Pythagorase:

$$b^2 = \left(\frac{|\langle \alpha|b\rangle|}{\alpha^2}\right)^2\alpha^2 + b_\perp^2 \ge \left(\frac{|\langle \alpha|b\rangle|}{\alpha^2}\right)^2\alpha^2 = \frac{|\langle \alpha|b\rangle|^2}{\alpha^2}$$

 $a^2b^2 \ge |\langle a|b\rangle|^2 \Rightarrow |a||b| \ge |\langle a|b\rangle| \stackrel{\text{lik}}{\ge} \langle a|b\rangle$

|| has two meanings: |complex number|

⇒ triangle inequality (in ℝ)

$$|a+b| \leq |a|+|b| \ \text{or} \ |x-z| \leq |x-y|+|y-z|$$

i.e., |a-b| is a **metric**.

Common shortcut: $a^2 \equiv |a|^2 = \langle a|a \rangle$

eIn complex spaces: $\langle b|\alpha\rangle^* = \langle a|b\rangle$ and for scalar $c \in \mathbb{C}$ it holds $|ca|^2 = \langle ca|ca\rangle = c^*c\langle a|\alpha\rangle = |c|^2|\alpha|^2$

Covector example

Example. Force $\vec{F} = \text{covector}$, path $d\vec{s} = \text{vector}$.

$$\vec{F} = -\vec{\nabla}U$$
, $dW = \vec{F} \cdot d\vec{s}$

Units: $[\vec{F}]$ = energy/length, $[d\vec{s}]$ = length.

If length unit changes from m to cm. d\$\vec{s}\$ multiplies 100×, but (if the energy unit remains the same) \vec{F} multiplies $0.01 \times$.

Maple

In package LinearAlgebra, operator "." is used for scalar product:

covector.vector rows.columns (in matrix multiplication)

+ = transposition

^* = Hermitean conjugate

Hilbert space

Hilbert space = linear space with a scalar product which is:

example of not-complete space: finite linear combinations of $\{(1,0,0,\ldots),(0,1,0,\ldots),(0,0,1,\ldots),\ldots\}$ igoplus complete (any Cauchy sequence converges in the (u, u) metric) usually also separable (it contains a countable dense subset

⇒ there is a countable basis) Loosely: "no vector is missing"

"it is not too big" or "there are no problems with using infinite sums"

Any finite vector space is a Hilbert space.

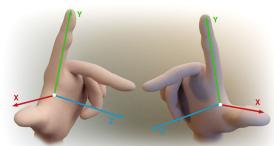
Example. Wavefunction is a vector of a Hilbert space, $\int |\psi(\tau)|^2 d\tau$ must be finite⁹. The scalar product is:

$$\langle \phi | \psi \rangle = \int \phi(\tau)^* \, \psi(\tau) \mathrm{d} \tau$$

n bosons: $\tau \in \mathbb{R}^{3n}$, *n* fermions (chemistry): $\tau \in (\mathbb{R} \times \{\alpha, \beta\})^{3n}$



3D: Right- and left-handed coordinate system



Left Handed Coordinates

Right Handed Coordinates

Right-handed: math, science, technology (Maple default) Left-handed: 3D image processing (Micro\$oft Direct 3D, PovRay)

^fSequence $\{v_i\}_{i=1}^{\infty}$ is Cauchy if $\forall d > 0 \ \exists n : |v_j - v_i| < d \ \forall i, j > n. \rightarrow \rightarrow$ 9For bound states , cf. de Broglie free-space "matter waves" .

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Unitary matrix Orthogonal^h (in \mathbb{R}^n) or **unitary** (in \mathbb{C}^n) matrix is a square matrix for which:

Square matrix $n \times n$, e.g.:

$$A = \begin{pmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{21} & A_{22} & A_{23} \end{pmatrix}$$

may represent:

igoplus matrix of coefficients of a set of n of linear equations for n unknowns:

$$\sum_{i} A_{ij} x_{j} = b_{i} \quad \text{or} \quad A \cdot x = b \quad \text{or} \quad Ax = b \quad \text{or} \quad |\hat{A}|x\rangle = |b\rangle$$

igoplus linear transformation (map, operator) $\mathbb{R}^n \to \mathbb{R}^n$ or $\mathbb{C}^n \to \mathbb{C}^n$

$$x_i \to \sum_i A_{ij} x_j$$
 or $x \to A \cdot x$ or $x \to Ax$ or $|x\rangle \to |\hat{A}|x\rangle$

igoplus matrix of coefficients of a quadratic form $\mathbb{R}^n \to \mathbb{R}$ or $\mathbb{C}^n \to \mathbb{C}$

$$x_i \to \sum_i x_i A_{ij} x_j \quad \text{or} \quad x \to x^\mathsf{T} \cdot A \cdot x \quad \text{or} \quad x \to x^\mathsf{T} A x \quad \text{or} \quad |x\rangle \to \langle x | \hat{A} | x \rangle$$

a quadratic tensor: e.g., of pressure or small deformation

$$\sum_{j} U^{\mathsf{T}}_{ij} U_{jk} = \sum_{j} U_{ji} U_{jk} = \delta_{ik} \ \text{or} \ \sum_{j} U^{\dagger}_{ij} U_{jk} = \sum_{j} U^{*}_{ji} U_{jk} = \delta_{ik}$$

 $U^{\mathsf{T}} \cdot U = \delta$ or $U^{\mathsf{T}} \cdot U = \delta$

- igcup columns U_{*i} can be treated as coordinates of an orthonormal basis (in other orthonormal basis), i.e., a (matrix of) unitary transformation
- \bigcirc *U* is regular: $U^{-1} = U^{\dagger}$

or in coordinates

- \bigcirc | det U| = 1 (in \mathbb{C}); in \mathbb{R} this means that det $U = \pm 1$
- a unitary matrix transforms an orthonormal basis to an orthonormal basis
- \bigcirc linear map $x \to U \cdot x$ "preserves angles", in $\mathbb R$ it can be interpreted as:
 - \bigcirc rotation in \mathbb{R}^n (for det U=1)
 - igoplus rotation and reflection in \mathbb{R}^n (for $\det U = -1$).

Examples of linear transformations in \mathbb{R}^n useful in molecular chemistry: mmpc1.mw

hterm "orthonormal" is not used

Matrices Notation:

In quantum theory often denoted as Â

- \bigcirc Other habits (e.g., as tensors): \overrightarrow{A} , \underline{A}
- \bigcirc $A \cdot x$ is less common than Ax; in the bra-ket notation $A|x\rangle$ or $|Ax\rangle$ or $|A|x\rangle$
- Vectors u and co-vectors u^{T} or $u^{\dagger} \equiv \langle u |$ ("bra") should be distinguished.

Matrices in infinite-dimension spaces are infinite = linear operators

If the set of equations $A \cdot x = b$ can be solved $\forall b$, then A is called **regular**. The solution is then:

 $x=A^{-1}\cdot b$

where $A^{-1} =$ **inverse matrix**, $A \cdot A^{-1} = A^{-1} \cdot A = \delta$, and $\delta =$ diag(1, 1, ...) =**unit matrix**, identity matrix, in coordinates Kronecker delta, also written as E, $\mathbf{1}$, I, I, I, etc.

Examples. Invert matrices:

Determinant

Determinant of a square matrix A is the number defined as a sum over all n! permutations p of indices $\{1, 2, \ldots, n\}$:

$$\det A = \sum_{p} \operatorname{sign}(p) \prod A_{i,p(i)}$$

where $sign(p) = (-1)^{number of transpositions in p}$.

 $\det A \neq 0$ for a regular matrix.

It holds

$$det(A \cdot B) = det(A) det(B), \quad det(A^{-1}) = \frac{1}{det A}$$
 (for regular A)

The determinant of a diagonal or triangular matrix = product of the numbers on the diagonal

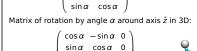
Example. Calculate a) sign(2, 3, 1), b) sign(n, n-1, n-2, ..., 2, 1)

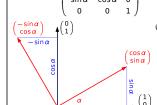
a) 1, b) (-1) $^{n(n-1)/2}$ (= 1 for $n \equiv 0$, 3 mod 4 and -1 otherwise)

Matrix of rotation

Matrix of rotation by oriented angle $+\alpha$ in 2D:

 $\cos \alpha - \sin \alpha$





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Internal coordinates:

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Write a matrix of rotation by angle α around vector $(a, b, c)^T$

Use spherical coordinates:

Example

 $(\alpha,b,c)=(r\sin\theta\cos\varphi,r\sin\theta\sin\varphi,r\cos\theta)$ reverse: $r = \sqrt{\alpha^2 + b^2 + c^2}$, $\theta = \arccos(c/r)$, $\varphi = \arctan(b, \alpha)$ Overloaded function $\arctan(b, a) = \arctan(b/a) + k\pi$. where k is such integer that $\varphi = \arctan(b, \alpha)$ is in the correct quadrant. In Fortran and C called atan2.

Compose from right (= in the order it is applied to a vector): = rotation by $-\varphi$ around \hat{z}

 R_1^{-1} = rotation by $-\varphi$ around \hat{z} R_2^{-1} = rotation by $-\theta$ around \hat{y}

 $R_3 = \text{rotation by } \alpha \text{ around } \hat{z}$

 R_2 = rotation by θ around \hat{y}

 $R_1 = \text{rotation by } \varphi \text{ around } \hat{z}$

Rotation matrix

see mmpc1.mw

 $R = R_1 \cdot R_2 \cdot R_3 \cdot R_2^{-1} \cdot R_1^{-1}$

φ