Vectors

Kindergarten: vector = $(v_1, ..., v_n)$, $v_i \in \mathbb{R}$ Quantum kindergarten: $v_i \in \mathbb{C}$

Mathematics: **vector space** (linear space) is defined by the axioms:

For vectors u, v, w and numbers $a, b \in \mathbb{R}$ or \mathbb{C} (a field^a in general):

u + (v + w) = (u + v) + wu + v = v + u $\exists \text{ null vector } 0: v + 0 = v$ $\exists \text{ opposite vector } -v: v + (-v) = 0$ a(bv) = (ab)v1v = va(u + v) = au + av(a + b)v = av + bv

} as in ℝ

Notation: v, \mathbf{v} , \vec{v} (real in 2D, 3D), \underline{v} , $|v\rangle$ ("ket"), v_i (?)

^ačesky komutativní těleso

Linear dependence

A set of nonzero vectors $v^{(i)}$, i = 1...m, is **linearly dependent** if there is a null linear combination with at least one of a_i nonzero:

A linearly independent set of vectors such that any vector (of given space) can be expressed as its linear combination is called a **basis**

 $v = \sum v_i b^{(i)}$

 $\sum a_i v^{(i)} = 0$

Example. Are the following vectors in \mathbb{R}^4 linearly dependent? $(1, 2, 3, 4)^{\top}, (1, -2, 3, -4)^{\top}, (1, 0, 1, 0)^{\top}$

Example. Are the following vectors in \mathbb{C}^2 linearly dependent?

 $cos(3x), \ldots$ Can function $cos^2(x)$ be expressed in this basis?

$$\left(\begin{array}{c}i\\1\end{array}\right),\left(\begin{array}{c}1+i\\1-i\end{array}\right)$$

see mmpc1.mw

ou

λG2

2/16

mmpc1

Scalar (inner, dot) product

We need a richer structure!

Kindergarten: scalar product $\vec{u} \cdot \vec{v} = \sum u_i v_i$

Mathematics: (u, v) is a number (real, complex) obeying axioms:

 $(u, v) = (v, u)^* (* = \text{complex conjugate})$ $(au, v) = a^*(u, v) \text{ (in physics)} \Rightarrow (u, av) = a(u, v)$ $= a(u, v) \text{ (in mathematics)} \Rightarrow (u, av) = a^*(u, v)$ (u + v, w) = (u, w) + (v, w) $(u, u) \ge 0$ $(u, u) = 0 \Rightarrow u = 0 \text{ (null vector)}$

Notation: $u^{\mathsf{T}}v$, $u^{\mathsf{T}}v$, $u^{\dagger}v$, \mathbf{b} (u, v), $\langle u, v \rangle$, $\mathbf{u} \cdot \mathbf{v}$, $\mathbf{u} \cdot \mathbf{v}$, $\langle u | v \rangle$ (bra-ket), $u_i v^i$ (covector-vector) \cdot (real spaces) or | (complex spaces) = sum over a pair of indices

Definition: If (u, v) = 0, vectors u, v are **perpendicular**

 $(u, u)^{1/2} = |u| = ||u|| = \text{norm}^{c}$

^bsymbol [†] = transpose + complex conjugate = adjoint = Hermitean (Hermitian) conjugate ^csimilar space with a norm only (and complete) = Banach space; under some conditions $(u, v) = (|u + v|^2 - |u - v|^2)/4$

(Cauchy-)Schwarz inequality

Dot-product in \mathbb{R}^n : $\vec{x} \cdot \vec{y} = |\vec{x}| |\vec{y}| \cos \theta \equiv xy \cos \theta \leq xy$. For nonzero *a*, *b* (zero cases are trivial):^d

$$b_{\perp} = b - \frac{\langle a|b \rangle}{a^2} a \implies \langle a|b_{\perp} \rangle = \langle a|b - \frac{\langle a|b \rangle}{a^2} a \rangle = \langle a|b \rangle - \frac{\langle a|b \rangle}{a^2} \langle a|a \rangle = 0$$
$$b = \frac{\langle a|b \rangle}{a^2} a + b_{\perp}$$

Pythagoras^e:

$$b^{2} = \left(\frac{|\langle a|b\rangle|}{a^{2}}\right)^{2} a^{2} + b_{\perp}^{2} \ge \left(\frac{|\langle a|b\rangle|}{a^{2}}\right)^{2} a^{2} = \frac{|\langle a|b\rangle|^{2}}{a^{2}}$$
$$a^{2} b^{2} \ge |\langle a|b\rangle|^{2} \Rightarrow |a||b| \ge |\langle a|b\rangle| \stackrel{\mathbb{R}}{\ge} \langle a|b\rangle$$

|| has two meanings: |complex number| and |vector|

 \Rightarrow triangle inequality (in \mathbb{R})

$$|a + b| \le |a| + |b|$$
 or $|x - z| \le |x - y| + |y - z|$

i.e., $|\alpha - b|$ is a **metric**.

^dCommon shortcut: $a^2 \equiv |a|^2 = \langle a|a \rangle$

•In complex spaces: $\langle b|a\rangle^* = \langle a|b\rangle$ and for scalar $c \in \mathbb{C}$ it holds $|ca|^2 = \langle ca|ca\rangle = c^*c\langle a|a\rangle = |c|^2|a|^2$

b⊥

(a|b)

h

a

Hilbert space

- Hilbert space = linear space with a scalar product which is:
- complete (any Cauchy sequence^f converges in the (u, u) metric)
- usually also separable (it contains a countable dense subset \Rightarrow there is a countable basis)
- Loosely: "no vector is missing"
 - "it is not too big" or "there are no problems with using infinite sums"

Any finite vector space is a Hilbert space.

Example. Wavefunction is a vector of a Hilbert space, $\int |\psi(\tau)|^2 d\tau$ must be finite⁹. The scalar product is:

$$\langle \phi | \psi \rangle = \int \phi(\tau)^* \psi(\tau) \mathrm{d} \tau$$

n bosons: $\tau \in \mathbb{R}^{3n}$, *n* fermions (chemistry): $\tau \in (\mathbb{R} \times \{\alpha, \beta\})^{3n}$

^fSequence $\{v_i\}_{i=1}^{\infty}$ is Cauchy if $\forall d > 0 \ \exists n : |v_j - v_i| < d \ \forall i, j > n. \rightarrow \rightarrow$ ^gFor bound states , cf. de Broglie free-space "matter waves" ... example of not-complete space: finite linear combinations of $\{(1, 0, 0, ...), (0, 1, 0, ...), ...\}$



Orthogonal and orthonormal bases

Orthogonal basis = all vectors are perpendicular. Orthonormal basis = also normalized.

$$b^{(i)} \cdot b^{(j)} = \delta_{ij}$$

Components of v in an orthonormal basis:

$$v_i = v \cdot b^{(i)} \Rightarrow v = \sum v_i b^{(i)} = (v_1, \ldots, v_n)_b$$

Scalar product:

$$u \cdot v = \sum u_i v_i$$

Scalar product in \mathbb{C} in physics

$$\langle u|v\rangle = \sum u_i^* v_i$$

Gram–Schmidt orthogonalization

A general basis $b^{(i)}$ can be orthogonalized by the Gram–Schmidt algorithm:

> $b^{(1)} := b^{(1)}/|b^{(1)}|$ $b^{(2)} := b^{(2)} - \langle b^{(1)}|b^{(2)}\rangle b^{(1)}, \ b^{(2)} := b^{(2)}/|b^{(2)}|$ $b^{(3)} := b^{(3)} - \langle b^{(1)}|b^{(3)}\rangle b^{(1)} - \langle b^{(2)}|b^{(3)}\rangle b^{(2)}, \ b^{(3)} := b^{(3)}/|b^{(3)}|$

":=" means "assign to" as in computer code.

Bases used in a Hilbert space are usually orthogonal or orthonormal

Example. Find all orthonormal bases $\{b^{(1)}, b^{(2)}\}$ in \mathbb{C}^2 for $b^{(1)} = (1, i)/\sqrt{2}$ $(b_1^{(1)} = 1, b_2^{(1)} = i)$

$$\left\langle \begin{pmatrix} 1 \\ i \end{pmatrix} \middle| \begin{pmatrix} x \\ y \end{pmatrix} \right\rangle = x - iy \stackrel{!}{=} 0 \implies x = iy \implies b^{(2)} = \frac{c}{\sqrt{2}} \left| \begin{pmatrix} i \\ 1 \end{pmatrix} \right\rangle, |c| = 1$$

Remember complex conjugate in the dot product, $i^* = -i$

more examples: see mmpc1.mw

Another notation: $\begin{vmatrix} 1 \\ i \end{vmatrix}' = \begin{pmatrix} 1 \\ -i \end{vmatrix}$





7/16

mmpc1

Linear forms

Linear form f (linear operator) assigns a number $f(v) \in \mathbb{R}$ (or \mathbb{C}) to a vector.

Axioms: for linear forms f, g, number a, and a vector v:

$$f(av) = f(v) + g(v)$$

$$f(av) = af(v)$$

For finite *n* one can write (In infinite-dimension spaces there may be continuity problems):

$$f(u) = \sum_{i=1}^{n} f_i u_i$$

Otherwise in Hilbert spaces linear form \approx scalar product:

$$f(\mathbf{v}) = \sum f_i \mathbf{v}_i = (f^*, \mathbf{v})$$

Linear form in Euclidean spaces (in some context) = covector, dual vector, covariant vector ("normal" vector = contravariant vector)

vector = column vector,

• covector = row vector (transposed) f^{T} , inverse transformation if a basis changes Scalar product then is: $f(u) = f^{\mathsf{T}} \cdot u = f^{\mathsf{T}} u = f^{i} u_{i}$ (Einstein summation convention). In complex Hilbert spaces $^{\mathsf{T}} \rightarrow ^{\dagger}$

Covector example



Example. Force \vec{F} = covector, path d \vec{s} = vector.

$$\vec{F} = -\vec{\nabla}U, \quad \mathrm{d}W = \vec{F}\cdot\mathrm{d}\vec{s}$$

Units: $[\vec{F}] = \text{energy/length}, [d\vec{s}] = \text{length}.$

If length unit changes from m to cm, $d\vec{s}$ multiplies $100 \times$, but (if the energy unit remains the same) \vec{F} multiplies $0.01 \times .$

Maple

In package LinearAlgebra, operator "." is used for scalar product:

covector.vector

rows.columns (in matrix multiplication)

^+ = transposition

^* = Hermitean conjugate

3D: Right- and left-handed coordinate system



credit: Wikipedia

Right-handed: math, science, technology (Maple default) Left-handed: 3D image processing (Micro\$oft Direct 3D, PovRay)

Matrices

Square matrix $n \times n$, e.g.:

$$A = \begin{pmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{pmatrix}$$

may represent:

matrix of coefficients of a set of *n* of linear equations for *n* unknowns:

$$\sum_{j} A_{ij} x_j = b_i \quad \text{or} \quad A \cdot x = b \quad \text{or} \quad Ax = b \quad \text{or} \quad |\hat{A}|x\rangle = |b\rangle$$

linear transformation (map, operator) $\mathbb{R}^n \to \mathbb{R}^n$ or $\mathbb{C}^n \to \mathbb{C}^n$

$$x_i \rightarrow \sum_j A_{ij} x_j$$
 or $x \rightarrow A \cdot x$ or $x \rightarrow Ax$ or $|x\rangle \rightarrow |\hat{A}|x\rangle$

matrix of coefficients of a quadratic form $\mathbb{R}^n \to \mathbb{R}$ or $\mathbb{C}^n \to \mathbb{C}$

$$x_i \rightarrow \sum_{ij} x_i A_{ij} x_j$$
 or $x \rightarrow x^T \cdot A \cdot x$ or $x \rightarrow x^T A x$ or $|x\rangle \rightarrow \langle x | \hat{A} | x \rangle$

a quadratic tensor; e.g., of pressure or small deformation

12/16 mmpc1

Matrices

Notation:

- In quantum theory often denoted as \hat{A}
- Other habits (e.g., as tensors): A, <u>A</u>
- $A \cdot x$ is less common than Ax; in the bra-ket notation $A|x\rangle$ or $|Ax\rangle$ or $|A|x\rangle$
- Vectors u and co-vectors u^{T} or $u^{\dagger} \equiv \langle u |$ ("bra") should be distinguished.
- Matrices in infinite-dimension spaces are infinite = linear operators
- If the set of equations $A \cdot x = b$ can be solved $\forall b$, then A is called **regular**. The solution is then:

$$x = A^{-1} \cdot b$$

where $A^{-1} =$ **inverse matrix**, $A \cdot A^{-1} = A^{-1} \cdot A = \delta$, and $\delta =$ diag(1, 1, ...) =**unit matrix**, identity matrix, in coordinates Kronecker delta, also written as E, 1, I, I, i, etc. **Examples.** Invert matrices:

a)
$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$
, b) $\begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$ $\begin{pmatrix} \tau & 0 \\ 0 & 1 \end{pmatrix}$ $\begin{pmatrix} \tau & 0 \\ \tau & \tau \end{pmatrix}$ (q ' $\begin{pmatrix} \epsilon/\tau & 0 & 0 \\ 0 & \tau/\tau & 0 \\ 0 & 0 & \tau \end{pmatrix}$ (e

Determinant

Determinant of a square matrix A is the number defined as a sum over all n! permutations p of indices $\{1, 2, ..., n\}$:

$$\det A = \sum_{p} \operatorname{sign}(p) \prod A_{i,p(i)}$$

where sign(p) = (-1)^{number of transpositions in p.}

det $A \neq 0$ for a regular matrix.

It holds

$$det(A \cdot B) = det(A) det(B), \quad det(A^{-1}) = \frac{1}{detA}$$
 (for regular A)

The determinant of a diagonal or triangular matrix = product of the numbers on the diagonal

Example. Calculate a) sign(2, 3, 1), b) sign(*n*, *n* – 1, *n* – 2, ..., 2, 1)

(a) 1, b) $(-1)^{n(n-1)/2}$ (= 1 for $n \equiv 0$, 3 mod 4 and -1 otherwise)

Unitary matrix

14/16 mmpc1

Orthogonal^h (in \mathbb{R}^n) or **unitary** (in \mathbb{C}^n) matrix is a square matrix for which:

$$U^{\mathsf{T}} \cdot U = \delta$$
 or $U^{\dagger} \cdot U = \delta$

or in coordinates

$$\sum_{j} U^{\mathsf{T}}_{ij} U_{jk} = \sum_{j} U_{ji} U_{jk} = \delta_{ik} \text{ or } \sum_{j} U^{\dagger}_{ij} U_{jk} = \sum_{j} U^{*}_{ji} U_{jk} = \delta_{ik}$$

- columns U_{*i} can be treated as coordinates of an orthonormal basis (in other orthonormal basis), i.e., a (matrix of) unitary transformation
- \bigcirc U is regular: $U^{-1} = U^{\dagger}$
- $|\det U| = 1$ (in \mathbb{C}); in \mathbb{R} this means that det $U = \pm 1$
- a unitary matrix transforms an orthonormal basis to an orthonormal basis
- linear map $x \to U \cdot x$ "preserves angles", in \mathbb{R} it can be interpreted as:
 - **)** rotation in \mathbb{R}^n (for det U = 1)
 - > rotation and reflection in \mathbb{R}^n (for det U = -1).

Examples of linear transformations in \mathbb{R}^n useful in molecular chemistry: mmpc1.mw ^hterm "orthonormal" is not used

Matrix of rotation

Matrix of rotation by oriented angle $+\alpha$ in 2D:

Matrix of rotation by angle α around axis \hat{z} in 3D:

 $\cos \alpha - \sin \alpha 0$

cosα





0

show pic/ATH 15/16 mmpc1

Internal coordinates:



Example

Write a matrix of rotation by angle α around vector $(a, b, c)^T$

Use spherical coordinates: $(a, b, c) = (r \sin \theta \cos \varphi, r \sin \theta \sin \varphi, r \cos \theta)$ reverse: $r = \sqrt{a^2 + b^2 + c^2}, \theta = \arccos(c/r), \varphi = \arctan(b, a)$ Overloaded function $\arctan(b, a) = \arctan(b/a) + k\pi$, where k is such integer that $\varphi = \arctan(b, a)$ is in the correct quadrant. In Fortran and C called atan2.

Compose from right (= in the order it is applied to a vector): R_1^{-1} = rotation by $-\varphi$ around \hat{z}

$$R_2^{-1}$$
 = rotation by $-\theta$ around \hat{y}

- R_3 = rotation by α around \hat{z}
- R_2 = rotation by θ around \hat{y}
- $R_1 =$ rotation by φ around \hat{z}
- Rotation matrix





see mmpc1.mw