

## Stlačitelné proudění, segregovaná a sdružená metoda

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### **Compressible flows & OF**

#### Segregated OF solvers for compressible flows:

SIMPLE based solvers: rhoSimpleFoam, rhoPimpleFoam

- different setup for subsonic and transonic/supersonic flows:
  - compressibility effects in the continuity equation
  - different relaxation factors (problem dependent?)
- very sensitive to initial condition
- low efficiency for transonic/supersonic flows

PISO based solvers: *sonicFoam*, hybrid AUSM+ scheme of [Xisto]

- limited Courant number  $\rightarrow$  limited efficiency for steady state cases

#### Coupled (density based) OF solvers:

- rhoCentralFoam: explicit central-upwind scheme by Tadmor
- *dbns* solvers (Foam-extend): explicit Riemann solvers based scheme
- aeroFoam: (?) similar to dbns
- **\*** All coupled OF schemes use explicit time stepping  $\rightarrow$  strong limit on  $\Delta t$ !



### Segregovaný řešič, algoritmus SIMPLE

#### Rovnice pro tlak z rovnice kontinuity

$$U = \hat{U} - \frac{1}{a} \nabla p$$
  

$$\rho = \psi p$$
  

$$\nabla \cdot (\rho U) = \nabla \cdot (\psi \hat{U} \mathbf{p}) - \nabla \cdot (\frac{\rho}{a} \nabla \mathbf{p}) = 0$$
  

$$\psi = \frac{\rho}{p} = \frac{1}{rT}$$
  

$$\psi = const. \Rightarrow T = const.$$

#### Jak odpovídá předpoklad *T=const.* rovnici energie?

$$H = c_p T + \frac{1}{2} ||U||^2 = const.$$
$$T' \approx -\frac{1}{c_p} U \cdot U'$$

$$\frac{T'}{T} \approx -\frac{U^2}{c_p T} \frac{U'}{U} = (1 - \gamma) M^2 \frac{U'}{U}$$



### Segregovaný řešič

#### **Eulerovy rovnice v 1D**

$$\rho_t + (\rho u)_x = 0$$
$$(\rho u)_t + (\rho u^2 + p)_x = 0$$
$$(\rho E)_t + (\rho u H)_x = 0$$

$$p = \rho r T$$
$$E = c_v T + u^2/2$$
$$H = c_p T + u^2/2$$

#### Linearizace (stacionární případ)

$$\rho'_{t} + \rho'_{x}\bar{u} + \bar{\rho}u'_{x} = 0 \qquad \frac{\rho'_{x}}{\bar{\rho}} + \frac{u'_{x}}{\bar{u}} = 0$$
$$u'_{t} + \bar{u}u'_{x} + \frac{1}{\bar{\rho}}p'_{x} = 0 \qquad \frac{u'_{x}}{\bar{u}} + \frac{\bar{p}}{\bar{\rho}\bar{u}^{2}}\frac{p'_{x}}{\bar{p}} = 0$$
$$p'_{t} + \bar{u}p'_{x} + \kappa\bar{p}u'_{x} = 0 \qquad \frac{u'_{x}}{\bar{u}} + \kappa\frac{p'_{x}}{\bar{p}} = 0$$

#### **Algoritmus SIMPLE**

$$\frac{u'}{\bar{u}} \approx -\frac{1}{a} \frac{\kappa}{M^2} \frac{p'_x}{\bar{p}},$$

$$\bar{u}\frac{p'_x}{\bar{p}} - \left(\frac{\kappa}{aM^2}\frac{p'_x}{\bar{p}}\right)_x = 0,$$

#### => relaxace závislá na M



## **LU-SGS Solver for OF**

#### **Previous works:**

- [Gill et al, 2013] LU-SGS for steady turbulent flows (at OFW8)
- [Heyns et al, 2014] GMRES/LU-SGS solver
- [Shen et al, 2016] detailed description of the implementation, dual time stepping for unsteady cases

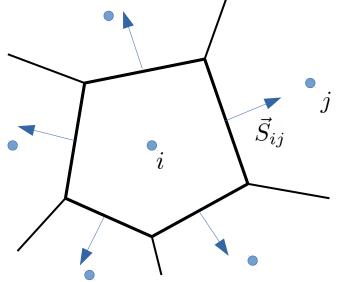
#### **Current work:**

- matrix-free LU-SGS
- based on *dbns* library with following improvements:
  - run-time selection of the Riemann solver (AUSM+up, HLLC, ...)
  - support for dynamic meshes (arbitrary Lagrangian-Eulerian method, multiple reference frames)
- unified solver for steady and transient case



## **LU-SGS Solver for OF**

- Navier-Stokes equations for compressible flows
- AUSM and HLLC fluxes
- central approximation of viscous fluxes
- low-order Jacobian based on Rusanov flux



$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u_j)}{\partial x_j} = 0$$

$$\frac{\partial (\rho u_i)}{\partial t} + \frac{\partial (\rho u_i u_j)}{\partial x_j} + \frac{\partial p}{\partial x_i} = \frac{\partial \tau_{ij}^{eff}}{\partial x_j}$$

$$\frac{\partial (\rho E)}{\partial t} + \frac{\partial \left[ (\rho E + p) u_j \right]}{\partial x_j} = \frac{\partial (u_i \tau_{ij}^{eff} - q_j^{eff})}{x_j}$$

$$W = [\rho, \rho \vec{u}, \rho E]$$
$$|\Omega_i| \frac{dW_i}{dt} = -R(W)_i = -\sum_j (F_{ij} - F_{ij}^v)$$

$$|\Omega_i| \frac{W_i^{n+1} - W_i^n}{\Delta t} \approx -R(W^n)_i - \frac{\partial R^{lo}}{\partial W_j} \left( W_j^{n+1} - W_j^n \right)$$

$$\sum_{j} \left[ \frac{|\Omega_i|}{\Delta t} \delta_{ij} + \frac{\partial R_i^{lo}}{\partial W_j} \right] \Delta W_j = -R(W^n)_i.$$



### **LU-SGS Solver for OF**

#### Lower-Upper Symmetric Gauss Seidel method for *Ax=b*:

 $A = L + D + U \approx (L + D)D^{-1}(D + U)$ 

$$D\Delta x^* = b - Ax^n - L\Delta x^*,$$
$$D\Delta x = D\Delta x^* - U\Delta x,$$
$$x^{n+1} = x^n + \Delta x.$$

#### Matrix-free LU-SGS for compressible flows [Blazek]:

$$D_i \Delta W_i^* = -\tilde{R}(W^n)_i - \frac{1}{2} \sum_{j < i} \left[ \Delta \mathbb{F}_j^* \cdot \vec{S}_{ij} + \lambda_{ij} \Delta W_j^* \right], \text{ for i=0,...,#cells,}$$
$$D_i \Delta W_i = D_i \Delta W_i^* - \frac{1}{2} \sum_{j > i} \left[ \Delta \mathbb{F}_j \cdot \vec{S}_{ij} + \lambda_{ij} \Delta W_j \right], \text{ for i=#cells,...,0.}$$

where:

$$\begin{split} \Delta \mathbb{F}_{j}^{*} &= \mathbb{F}(\vec{W}_{j}^{*}) - \mathbb{F}(\vec{W}_{j}^{n}), \\ \Delta \mathbb{F}_{j} &= \mathbb{F}(\vec{W}_{j}^{n+1}) - \mathbb{F}(\vec{W}_{j}^{n}), \end{split} \quad \mathbb{F}(W) = \begin{bmatrix} \rho \vec{u} \\ \rho \vec{u} \otimes \vec{u} + p \mathbb{I} \\ (\rho E + p) \vec{u} \end{bmatrix}, \ D_{i} = \left( \frac{|\Omega_{i}|}{\Delta \tau_{i}} + \frac{1}{2} \sum_{j \in N_{i}} \lambda_{ij} \right) I, \\ \lambda_{ij} &= \omega \left[ |\vec{U}_{ij} \cdot \vec{S}_{ij}| + |\vec{S}_{ij}|a_{ij} + \frac{|\vec{S}_{ij}|}{||\vec{x}_{i} - \vec{x}_{j}||} \max \left( \frac{4}{3\rho_{ij}}, \frac{\gamma}{\rho_{ij}} \right) \left( \frac{\mu}{Pr} + \frac{\mu_{T}}{Pr_{T}} \right) \right]. \end{split}$$



### **Boundary conditions**

#### "Standard" OF boundary conditions at subsonic inlet:

- $p_{tot}$  implemented as totalPressure  $\rightarrow p_f = f(U_f, p_{tot})$
- $T_{tot}$  implemented as totalTemperature  $\rightarrow T_f = f(U_f, T_{tot})$
- $\alpha$  implemented as pressureDirectedVelocity  $\rightarrow$  U<sub>f</sub> = f( $\varphi_{f}, \alpha$ )

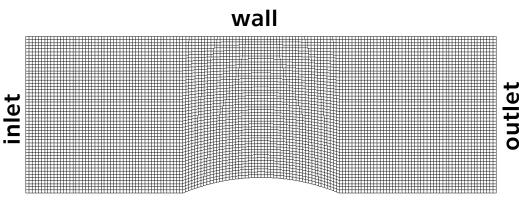
OK for segregated solvers where  $\varphi$  is updated from the flux of *pEqn*, but not compatible with coupled solvers where  $\varphi = f(U)$ !

#### New boundary condition for U at subsonic inlet:

1) 
$$U_f = f(p_{int'} p_{tot'} T_{tot'} \alpha),$$
  
2)  $p_f = f(U_f, p_{tot}),$  (totalPressure)  
3)  $T_f = f(U_f, T_{tot}),$  (totalTemperature)  
4)  $\varphi_f = f(U_f).$ 



### 2D flow over a bump



#### **Channel with circular bump**

length: 3m height: 1m bump height:

0.1m for subsonic inlets 0.04m for supersonic inlets

wall

#### **Test cases:** subsonic: $M_{in} = 0.1$ transonic: $M_{in} = 0.675$ supersonic: $M_{in} = 1.65$

#### Mesh:

structured coarse with 150x50 cells fine with 450x150 cells

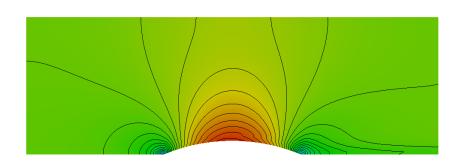
#### Flow field: 2D compressible ideal gas inviscid

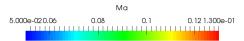
#### Solution obtained with

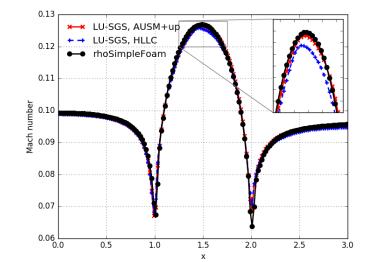
pressure correction method (rhoSimpleFOAM) LU-SGS with AUSM+up or HLLC scheme using Barth or Venkatakrishnan limiter

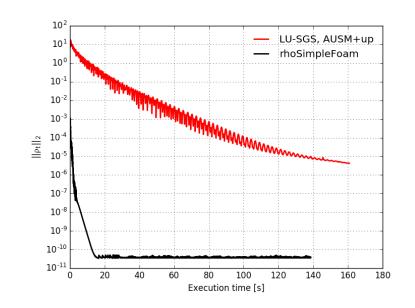


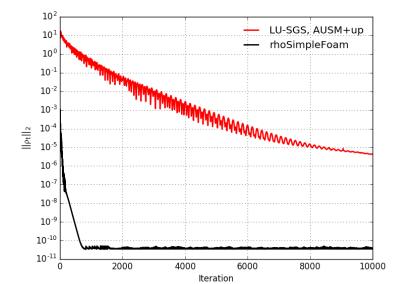






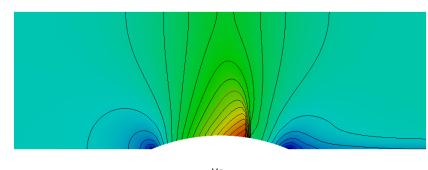




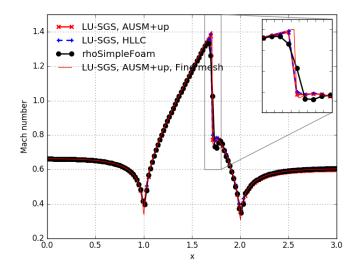


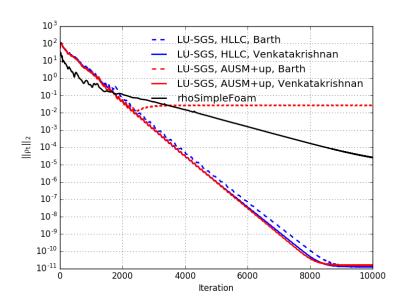
# Flow o. a bump, $M_{in}$ =0.675

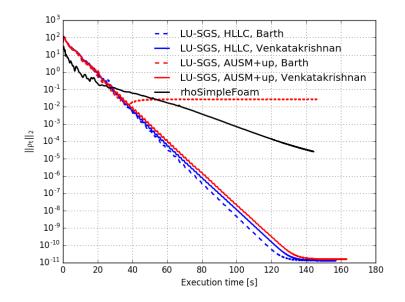




Ma 4.000e-01 0.75 1 1.25 1.400e+00

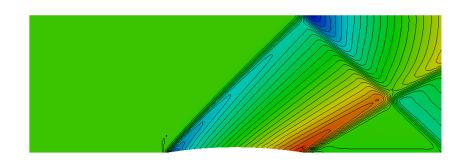




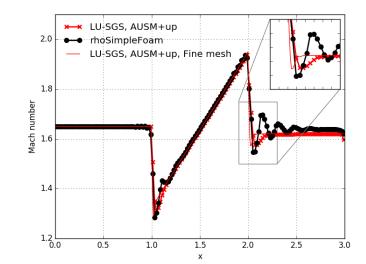


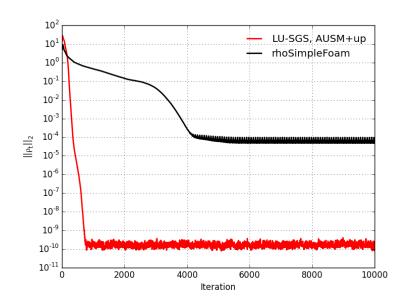
# Flow o. a bump, $M_{in}$ =1.65

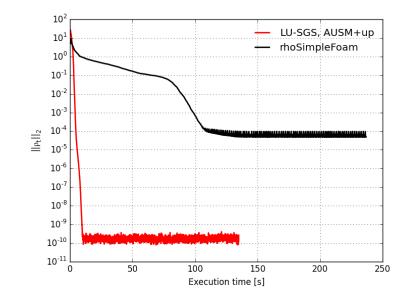




Ma 1.250e+00 1.4 1.57 1.75 1.950e+00





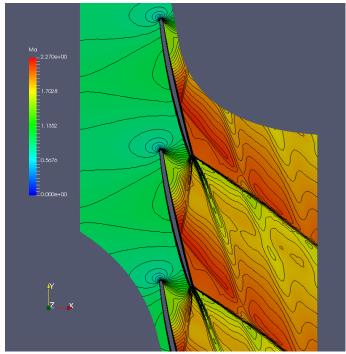




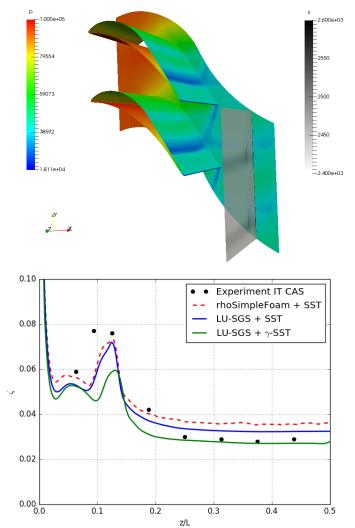
### **Applications**

### 2D flows through tip section of a turbine cascade

- M<sub>out</sub> ~ 2
- Re ~ 10<sup>6</sup>



#### 3D flows through turbine cascade

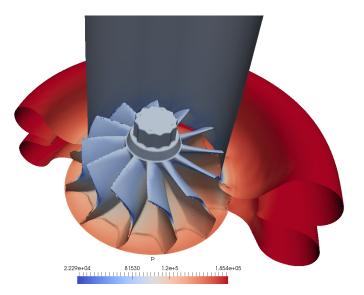


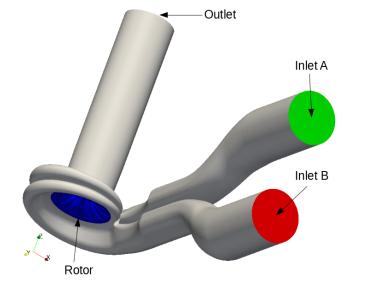


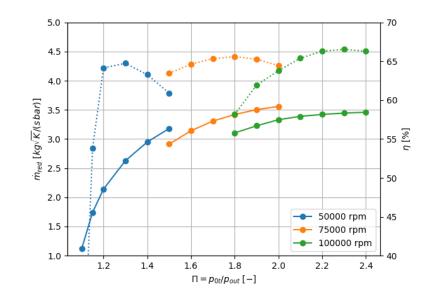
## **Applications**

#### Analysis of twin-scroll turbine

- high rotation speed (50-100 krpm)
- expansion ratio Π=1.2-2.2
- important compressibility
- steady (MRF) / transient simulation
- 2 mil. cells, snappyHexMesh









### Conclusions

The matrix-free LU-SGS method has been implemented for steady and unsteady turbulent flows (using MRF or dynamic mesh).

The LU-SGS method is:

- very simple
- efficient for compressible flows (Ma>0.3)
- provides sharp resolution of shock waves
- has very low memory footprint

The LU-SGS solver is compatible with the rest of OpenFOAM framework (turbulence models, parallel processing, ...)

#### Future work:

• preconditioning for low Mach number flows

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